

Reversibility conditions for quantum operations

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Some basic references

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Introduction

A quantum system is represented by a C^* -algebra $\mathcal{A} \subseteq B(\mathcal{H})$:

- Observables = self-adjoint operators: \mathcal{A}^h
- States of the system = positive unital functionals: $\mathcal{S}(\mathcal{A})$

Suppose another system is represented by another algebra \mathcal{B} , with state space $\mathcal{S}(\mathcal{B})$.

A **quantum operation** is a map $T : \mathcal{S}(\mathcal{A}) \rightarrow \mathcal{S}(\mathcal{B})$ representing a physical transformation of the states.

Introduction

Suppose $\mathcal{S} \subset \mathcal{S}(\mathcal{A})$, $T : \mathcal{S}(\mathcal{A}) \rightarrow \mathcal{S}(\mathcal{B})$ a quantum operation.

\mathcal{S} can be seen as carrying some information, this is lost under T :

- in quantum statistics, \mathcal{S} represents a statistical model, T a stochastic map (e.g. a restriction to a subalgebra)
- in error correction, \mathcal{S} represents a code, T a noisy channel

Reversibility/sufficiency: the set \mathcal{S} can be completely recovered by a quantum operation $S : \mathcal{S}(\mathcal{B}) \rightarrow \mathcal{S}(\mathcal{A})$.

The setting

- Let \mathcal{A} be a finite dimensional C^* -algebra,

$$\mathcal{A} \equiv \bigoplus_{k=1}^n B(\mathcal{H}_k) \subseteq B(\mathcal{H}), \quad \dim(\mathcal{H}) < \infty$$

Let Tr be the trace inherited from $B(\mathcal{H})$.

- \mathcal{A} is a (finite dim) Hilbert space, with the Hilbert-Schmidt inner product

$$\langle a, b \rangle = \text{Tr } a^* b, \quad a, b \in \mathcal{A}$$

- State space = density operators:

$$\mathcal{S}(\mathcal{A}) = \{\rho \geq 0, \text{Tr } \rho = 1\}$$

Quantum operations

- Quantum operations = completely positive trace preserving maps (channels) $T : \mathcal{A} \rightarrow \mathcal{B}$:

$$T(a) = \sum_i V_i^* a V_i, \quad a \in \mathcal{A}, \quad \sum_i V_i V_i^* = I_{\mathcal{A}}$$

- The adjoint $T^* : \mathcal{B} \rightarrow \mathcal{A}$ with respect to the HS inner product,

$$\langle a, T^*(b) \rangle = \langle T(a), b \rangle, \quad a, b \in \mathcal{A}$$

is completely positive and unital, $T^*(I_{\mathcal{B}}) = I_{\mathcal{A}}$.

Reversibility (sufficiency)

Let $T : \mathcal{A} \rightarrow \mathcal{B}$ be a quantum operation and let $\mathcal{S} \subset \mathcal{S}(\mathcal{A})$ be a set of states.

Definition

We say that T is **reversible** (or **sufficient**) for \mathcal{S} if there is a quantum operation $S : \mathcal{B} \rightarrow \mathcal{A}$, such that

$$S \circ T(\sigma) = \sigma, \quad \sigma \in \mathcal{S}$$

Assumptions

We will suppose:

1. $T : \mathcal{A} \rightarrow \mathcal{B}$ is a quantum operation
2. $\mathcal{S} \subset \mathcal{S}(\mathcal{A})$ and there is an invertible element $\rho \in \mathcal{S}$, such that $T(\rho)$ is invertible.

Let (\mathcal{S}, T) be such that $\mathcal{S} \subset \mathcal{S}(\mathcal{A})$ and T a quantum operation.
Then there is a projection $p \in \mathcal{A}$, such that

- $\mathcal{S}' = \text{co}(\mathcal{S}) \subseteq \mathcal{S}(p\mathcal{A}p)$,
- with $T' = T|_{p\mathcal{A}p}$, (\mathcal{S}', T') satisfy 1. and 2.
- T is reversible for \mathcal{S} if and only if T' is reversible for \mathcal{S}' .

The dual map

Let $T : \mathcal{A} \rightarrow \mathcal{B}$ be a quantum operation and $\rho \in \mathcal{S}(\mathcal{A})$ be such that ρ and $T(\rho)$ are invertible. The map $T_\rho^* : \mathcal{B} \rightarrow \mathcal{A}$

$$T_\rho^*(b) := \rho^{1/2} T^*(T(\rho)^{-1/2} b T(\rho)^{-1/2}) \rho^{1/2}, \quad b \in \mathcal{B}$$

is called the dual of T with respect to ρ .

- T_ρ^* is a quantum operation
- $T_\rho^* \circ T(\rho) = \rho$

Relative entropy

Let $\sigma, \rho \in \mathcal{S}(\mathcal{A})$. Relative entropy is defined by

$$S(\sigma, \rho) = \begin{cases} \text{Tr } \sigma(\log \sigma - \log \rho) & \text{if } \text{supp } \sigma \leq \text{supp } \rho \\ \infty & \text{otherwise} \end{cases}$$

- **Monotonicity** If $T : \mathcal{A} \rightarrow \mathcal{B}$ is a quantum operation, then

$$S(\sigma, \rho) \geq S(T(\sigma), T(\rho))$$

- If T is reversible, then

$$S(\sigma, \rho) = S(T(\sigma), T(\rho)), \quad \sigma \in \mathcal{S}$$

A theorem of Petz (1986)

Theorem

Suppose both ρ and $T(\rho)$ are invertible. TFAE:

- (i) T is reversible for $\{\sigma, \rho\}$.
- (ii) $S(\sigma, \rho) = S(T(\sigma), T(\rho))$.
- (iii) $T_\rho^* \circ T(\sigma) = \sigma$.

The dual map T_ρ^* plays the role of the recovery map S .

Corollary

Under assumptions 1. and 2., T is reversible for \mathcal{S} if and only if T is reversible for all pairs $\{\sigma_1, \sigma_2\} \subset \mathcal{S}$.

Conditions for reversibility

There are several conditions characterizing reversibility:

1. Quantum f -divergences (relative entropy, Rényi relative entropy,...)
2. Radon-Nikodym derivative
3. Factorization of states in \mathcal{S} (quantum factorization criterion)
4. Distinguishability measures obtained from hypothesis testing:
 L_1 - distance, quantum Chernoff bound, quantum Hoeffding bound
5. Quantum Fisher information and quantum χ^2 -distance

Quantum f -divergence

Let $\sigma, \rho \in \mathcal{S}(\mathcal{A})$, ρ is invertible.

- \mathcal{A} is a Hilbert space, with the HS inner product
- Relative modular operator: $\Delta_{\sigma, \rho} : \mathcal{A} \rightarrow \mathcal{A}$, positive,

$$\Delta_{\sigma, \rho}(a) = \sigma a \rho^{-1}$$

- Let $f : \mathbb{R}^+ \rightarrow \mathbb{R}$ be a function
- Quantum f -divergence:

$$S_f(\sigma, \rho) = \langle \rho^{1/2}, f(\Delta_{\sigma, \rho})(\rho^{1/2}) \rangle$$

Operator convex functions

- A function $f : \mathbb{R}^+ \rightarrow \mathbb{R}$ is **operator convex** if

$$f(tA + (1 - t)B) \leq tf(A) + (1 - t)f(B)$$

for $t \in [0, 1]$, $A, B \in B(\mathcal{H})^+$, $\dim(\mathcal{H}) < \infty$

- Integral representation (Hiai, Mosonyi, Petz, Beny 2011):

$$f(x) = f(0) + \alpha x + \beta x^2 + \int_{(0, \infty)} \left(\frac{x}{1+t} - \frac{x}{x+t} \right) d\mu_f(t),$$

$\alpha \in \mathbb{R}$, $\beta \geq 0$ and μ_f is a non-negative measure on $(0, \infty)$ satisfying $\int (1+t)^{-2} d\mu_f(t) < \infty$.

Monotonicity and convexity

If f is operator convex, then S_f is **non-increasing** under quantum operations:

$$S_f(\sigma, \rho) \geq S_f(T(\sigma), T(\rho))$$

- S_f is **jointly convex**:

$$S_f\left(\sum_i p_i \sigma_i, \sum_i p_i \rho_i\right) \leq \sum_i p_i S_f(\sigma_i, \rho_i)$$

where $p_i \geq 0$, $\sum p_i = 1$, $\sigma_i, \rho_i \in \mathcal{S}(\mathcal{A})$.

- $S_f(\sigma, \rho) \geq f(1)$

S_f is a distance (distinguishability) measure on $\mathcal{S}(\mathcal{A})$.

Examples

- Relative entropy,

$$f(x) = x \log x = \int_0^\infty \left(\frac{x}{1+t} - \frac{x}{x+t} \right) dt$$

- For $\alpha \in [0, 1]$, $S_\alpha(\sigma, \rho) = 1 - \text{Tr} \sigma^\alpha \rho^{1-\alpha}$, here

$$f_\alpha(x) = 1 - x^\alpha = 1 - x + \int_0^\infty \left(\frac{x}{1+t} - \frac{x}{x+t} \right) d\mu_\alpha(t)$$

$$d\mu_\alpha(t) = \frac{\sin(\alpha\pi)}{\pi} t^{\alpha-1} dt.$$

- Rényi relative entropy with $\alpha \in (0, 1)$,

$$R_\alpha(\sigma, \rho) = \frac{1}{\alpha - 1} \log(1 - S_\alpha(\sigma, \rho))$$

- $S_f(\sigma, \rho) = \text{Tr} \rho(L_\sigma + R_\rho)^{-1}(\rho)$, here

$$f(x) = \frac{1}{1+x}, \quad \mu_f = \delta_{\{1\}}$$

Reversibility

Theorem

Under the conditions 1. and 2., TFAE:

- (i) T is reversible for \mathcal{S} .
- (ii) $S_f(\sigma, \rho) = S_f(T(\sigma), T(\rho))$ for all $\sigma \in \mathcal{S}$ and all operator convex f
- (iii) (Hiai, Mosonyi, Petz, Bény, 2011)
 $S_f(\sigma, \rho) = S_f(T(\sigma), T(\rho))$ for all $\sigma \in \mathcal{S}$ and some operator convex f with $|\text{supp } \mu_f| \geq \dim(\mathcal{H})^2 + \dim(\mathcal{K})^2$.
- (iv) $T_\rho^* \circ T(\sigma) = \sigma$ for all $\sigma \in \mathcal{S}$.

An example and a counterexample

- Since $\text{supp } \mu_\alpha = (0, \infty)$ for $\alpha \in (0, 1)$, T is reversible for \mathcal{S} if and only if

$$\text{Tr } \sigma^\alpha \rho^{1-\alpha} = \text{Tr } T(\sigma)^\alpha T(\rho)^{1-\alpha}, \quad \sigma \in \mathcal{S}$$

for some $\alpha \in (0, 1)$ (Petz 1988, Petz & AJ 2006).

- It can be shown that the equality

$$S_f(\sigma, \rho) = S_f(T(\sigma), T(\rho)), \quad \sigma \in \mathcal{S}$$

for the function

$$f(x) = \frac{1}{1+x}, \quad \text{supp } \mu_f = \{1\}$$

does not imply reversibility of T .

A quantum Radon-Nikodym derivative

Let $\sigma, \rho \in \mathcal{S}(\mathcal{A})$, ρ invertible.

- The (commutant) Radon-Nikodym derivative

$$d = d(\sigma, \rho) := \rho^{-1/2} \sigma \rho^{-1/2}$$

- $d \geq 0$, $\|d\| = \inf\{\lambda > 0, \sigma \leq \lambda\rho\}$, $\|d\| = 1$ iff $\sigma = \rho$.
- For $T : \mathcal{A} \rightarrow \mathcal{B}$ a quantum operation, let $T_\rho := (T_\rho^*)^* : \mathcal{A} \rightarrow \mathcal{B}$ be the adjoint of T_ρ^* , so that

$$T_\rho(a) = T(\rho)^{-1/2} T(\rho^{1/2} a \rho^{1/2}) T(\rho)^{-1/2}.$$

Then

$$T_\rho(d(\sigma, \rho)) = d(T(\sigma), T(\rho))$$

Reversibility

Theorem

Under the conditions 1. and 2., let $T^* : \mathcal{B} \rightarrow \mathcal{A}$ be the adjoint of T . TFAE:

- (i) T is reversible for \mathcal{S} .
- (ii) $T^*(d(T(\sigma), T(\rho))) = d(\sigma, \rho)$, $\sigma \in \mathcal{S}$.
- (iii) $\Phi(d(\sigma, \rho)) = d(\sigma, \rho)$, $\sigma \in \mathcal{S}$, where $\Phi = T^* \circ T_\rho$.

Note that $\Phi^* = T_\rho^* \circ T : \mathcal{A} \rightarrow \mathcal{A}$ is a quantum operation, such that $\Phi^*(\rho) = \rho$. Moreover, (iii) means that $d(\sigma, \rho)$ is a fixed point of Φ .

The fixed point algebra

Let $\Phi : \mathcal{A} \rightarrow \mathcal{A}$ be a map such that Φ^* is a quantum operation. Suppose that $\Phi^*(\rho) = \rho$ for some invertible $\rho \in \mathcal{S}(\mathcal{A})$.

Denote $\mathcal{F}_\Phi = \{a \in \mathcal{A}, \Phi(a) = a\}$.

- \mathcal{F}_Φ is a subalgebra in \mathcal{A} .
- $\rho^{it} \mathcal{F}_\Phi \rho^{-it} \subseteq \mathcal{F}_\Phi$, $t \in \mathbb{R}$.
- ρ can be decomposed as

$$\rho = \rho^A \rho^B$$

where $\rho^A \in \mathcal{F}_\Phi^+$ and $\rho^B \in (\mathcal{F}'_\Phi \cap \mathcal{A})^+$.

An extension of \mathcal{S}

- Let T be reversible for \mathcal{S} , $\sigma \in \mathcal{S}$, $d = d(\sigma, \rho)$.
- Then $d \in \mathcal{F}_\Phi$.
- Hence if $\sigma_t := \rho^{it}\sigma\rho^{-it}$, then

$$d_t := d(\sigma_t, \rho) = \rho^{it}d\rho^{-it} \in \mathcal{F}_\Phi.$$

Corollary

T is reversible for \mathcal{S} if and only if T is reversible for $\{\rho^{it}\mathcal{S}\rho^{-it}, t \in \mathbb{R}\}$.

Factorization of states in \mathcal{S}

Theorem

Under the conditions 1. and 2., TFAE:

- (i) *T is reversible for \mathcal{S} .*
- (ii) *There is some $\rho^B \in (\mathcal{F}'_\Phi \cap \mathcal{A})^+$, such that for $\sigma \in \mathcal{S}$,*

$$\sigma = \sigma^A \rho^B$$

with $\sigma^A \in \mathcal{F}_\Phi$.

- (iii) *There is some $\rho^B \in \mathcal{A}^+$, such that for all $\sigma \in \mathcal{S}$,*

$$\sigma = T^*(\sigma_0^A) \rho^B, \quad T(\sigma) = \sigma_0^A T(\rho^B),$$

where $\sigma_0^A \in \mathcal{B}^+$.

Quantum hypothesis testing

Let $\sigma, \rho \in \mathcal{S}(\mathcal{A})$, ρ invertible. Let $H_0 = \rho$, $H_1 = \sigma$.

- Any test of H_0 against H_1 is given by some $0 \leq M \leq I$ in \mathcal{A} .
- M corresponds to rejecting the H_0 hypothesis.
- error of 1. kind:
rejecting true - probability $\alpha(M) = \text{Tr } M\rho$.
- error of 2. kind:
not rejecting false - probability $\beta(M) = \text{Tr } (I - M)\sigma$.
- For some $\lambda \in [0, 1]$, minimize

$$\lambda\alpha(M) + (1 - \lambda)\beta(M)$$

to obtain the **Bayes optimal test**.

Quantum Neyman-Pearson lemma (Helstrom,Holevo)

Let $\lambda \in [0, 1]$. The **minimum Bayes error probability** is

$$\begin{aligned}\Pi_\lambda &:= \inf_{0 \leq M \leq I} (\lambda \alpha(M) + (1 - \lambda) \beta(M)) \\ &= \frac{1}{2} (1 - \|(1 - \lambda)\sigma - \lambda\rho\|_1),\end{aligned}$$

where $\|a\|_1 = \text{Tr} |a| = \text{Tr} (a^* a)^{1/2}$ is the L_1 -norm. It satisfies monotonicity

$$\|a\|_1 \geq \|T(a)\|_1, \quad a \in \mathcal{A}$$

if T is a quantum operation.

Reversibility and L_1 -norm

Let $\tilde{\mathcal{S}} = \bigcup\{\rho^{is}\mathcal{S}\rho^{-is}, s \in \mathbb{R}\}$ be the extended family.

Theorem

Suppose 1. and 2. hold. Then

(i) T is reversible for \mathcal{S} if and only if

$$\|\sigma - t\rho\|_1 = \|T(\sigma) - tT(\rho)\|_1, \quad \sigma \in \tilde{\mathcal{S}}, t \geq 0$$

(ii) If \mathcal{B} is abelian, then T is reversible for \mathcal{S} if and only if

$$\|\sigma - t\rho\|_1 = \|T(\sigma) - tT(\rho)\|_1, \quad \sigma \in \mathcal{S}, t \geq 0$$

Moreover, in this case, all elements in \mathcal{S} commute.

Hypothesis testing for n copies

Let $\sigma, \rho \in \mathcal{S}(\mathcal{A})$

- Let $n \in \mathbb{N}$, n copies $\sigma^{\otimes n}, \rho^{\otimes n} \in \mathcal{S}(\mathcal{A}^{\otimes n})$
- $H_0 = \rho^{\otimes n}$, $H_1 = \sigma^{\otimes n}$
- A test $0 \leq M_n \leq I$, $M_n \in \mathcal{A}^{\otimes n}$
- Minimum Bayes error probability

$$\Pi_{\lambda,n} = \frac{1}{2}(1 - \|(1 - \lambda)\sigma^{\otimes n} - \lambda\rho^{\otimes n}\|_1)$$

- As $n \rightarrow \infty$, $\Pi_{\lambda,n} \rightarrow 0$ exponentially fast

Quantum Chernoff distance

Theorem

(Audenaert et al. (2008)) For all $\lambda \in [0, 1]$.

$$\lim_n -\frac{1}{n} \log \Pi_{\lambda,n} = -\log \left(\inf_{0 \leq s \leq 1} \text{Tr} \sigma^s \rho^{1-s} \right) =: \xi_{QCB}(\sigma, \rho)$$

ξ_{QCB} is called the **quantum Chernoff distance**. Properties:

- **non-increasing** under quantum operations:

$$\xi_{QCB}(\sigma, \rho) \geq \xi_{QCB}(T(\sigma), T(\rho))$$

- $\xi_{QCB}(\sigma, \rho) = 0$ if and only if $\sigma = \rho$.
- not a quantum f -divergence

QCB and reversibility

Suppose $\sigma, \rho \in \mathcal{S}(\mathcal{A})$, ρ and $T(\rho)$ invertible.

- There is some $s_0 \in [0, 1]$ such that

$$\inf \operatorname{Tr} T(\sigma)^s T(\rho)^{1-s} = \operatorname{Tr} T(\sigma)^{s_0} T(\rho)^{1-s_0}$$

- We have

$$\begin{aligned} \xi_{QCB}(T(\sigma), T(\rho)) &= -\log \operatorname{Tr} T(\sigma)^{s_0} T(\rho)^{1-s_0} \\ &\leq -\log \operatorname{Tr} \sigma^{s_0} \rho^{1-s_0} \leq \xi_{QCB}(\sigma, \rho) \end{aligned}$$

- Suppose $\xi_{QCB}(T(\sigma), T(\rho)) = \xi_{QCB}(\sigma, \rho)$, then

$$\operatorname{Tr} \sigma^{s_0} \rho^{1-s_0} = \operatorname{Tr} T(\sigma)^{s_0} T(\rho)^{1-s_0}$$

QCB and reversibility

- Suppose $s_0 \in (0, 1)$. Then the equality implies that T is reversible for $\{\sigma, \rho\}$.
- Suppose $s_0 = 0, 1$. Then if $T(\sigma)$ is invertible, $\xi_{QCB}(\sigma, \rho) = \xi_{QCB}(T(\sigma), T(\rho)) = 0$, hence $\sigma = \rho$.
- If σ is invertible, then T is reversible for $\{\sigma, \rho\}$ iff

$$\xi_{QCB}(\sigma, \rho) = \xi_{QCB}(T(\sigma), T(\rho)).$$

- There are examples with σ not invertible, where the equality holds while T is not reversible.

QCB and reversibility

Theorem

Under the conditions 1. and 2., TFAE:

- (i) *T is reversible for \mathcal{S} .*
- (ii) *$\xi_{QCB}(T(\sigma), T(\rho)) = \xi_{QCB}(\sigma, \rho)$, for all $\sigma \in \text{co}(\mathcal{S})$.*
- (iii) *$\|T^{\otimes n}(\sigma^{\otimes n}) - tT(\rho^{\otimes n})\|_1 = \|\sigma^{\otimes n} - t\rho^{\otimes n}\|_1$, $\sigma \in \mathcal{S}$, $t \geq 0$, $n \in \mathbb{N}$.*

If all elements in \mathcal{S} are invertible, we may replace $\text{co}(\mathcal{S})$ by \mathcal{S} in (ii).

Quantum Hoeffding distance

Asymmetric setting:

- The error probability $\alpha(M)$ is bounded, $\alpha(M) \leq \epsilon$
- The error probability $\beta(M)$ is optimized under this constraint:

$$\beta_\epsilon := \inf\{\beta(M), 0 \leq M \leq I, \alpha(M) \leq \epsilon\}$$

n copies, $M_n \in \mathcal{A}^{\otimes n}$ a sequence of tests of $\rho^{\otimes n}$ against $\sigma^{\otimes n}$.

- We require that $\alpha(M_n)$ decays exponentially fast:
 $\alpha(M_n) \leq e^{-nr}$ for some $r > 0$.
- Let

$$\beta_{n,r} := \inf\{\beta(M_n), 0 \leq M_n \leq I, \alpha(M_n) \leq e^{-nr}\}$$

Quantum Hoeffding distance

Theorem

(Audenaert et al. (2008))

$$\lim_n -\frac{1}{n} \log \beta_{n,r} = \sup_{0 \leq s < 1} \frac{-sr - \log \operatorname{Tr} \sigma^s \rho^{1-s}}{1-s} =: H_r(\sigma, \rho)$$

$H_r(\sigma, \rho)$ is the **quantum Hoeffding distance**. Suppose ρ and σ are both invertible. Then

- $H_r(\sigma, \rho) \geq H_r(T(\sigma), T(\rho))$ for a quantum operation T
- $H_0(\sigma, \rho) := \lim_{r \rightarrow 0+} H_r(\rho, \sigma) = S(\sigma, \rho)$
- The function $r \mapsto H_r(\sigma, \rho)$ is strictly convex and decreasing for $r \in [0, S(\rho, \sigma)]$
- $H_r(\sigma, \rho) = 0$ for all $r \geq S(\rho, \sigma)$

QHD and reversibility

If σ , ρ and $T(\rho)$ are invertible, then T is reversible for $\{\sigma, \rho\}$ iff $H_r(\sigma, \rho) = H_r(T(\sigma), T(\rho))$ for some $r \in [0, S(T(\rho), T(\sigma))]$.

Theorem

Suppose the conditions 1. and 2. hold. Let $X \subset \mathcal{S}$ be such that $\mathcal{S} \subseteq \overline{\text{co}}(X \cup \{\rho\})$ and $T(\rho) \notin \overline{T(X)}$. Then there is some $r_0 > 0$, such that TFAE:

- (i) T is reversible for \mathcal{S} .
- (ii) $H_r(\sigma, \rho) = H_r(T(\sigma), T(\rho))$ for all $\sigma \in \mathcal{S}$ and $r \geq 0$
- (iii) $H_r(\sigma, \rho) = H_r(T(\sigma), T(\rho))$ for all $\sigma \in \text{co}(\mathcal{S})$ and some $r \in [0, r_0]$.

If all $\sigma \in \mathcal{S}$ are invertible, we may replace $\text{co}(\mathcal{S})$ by X in (iii) and $r_0 = \inf_{\sigma \in X} S(T(\rho), T(\sigma))$.

Quantum Fisher information

Let $\mathcal{D} = \mathcal{D}(\mathcal{A})$ be the set of invertible elements in $\mathcal{S}(\mathcal{A})$.

- \mathcal{D} is a differentiable manifold.
- Tangent space at $\rho \in \mathcal{D}$: $\mathcal{T}_\rho = \{x = x^* \in \mathcal{A}, \text{Tr } x = 0\}$
- A Riemannian metric in \mathcal{D} : g_ρ is an inner product in \mathcal{T}_ρ , $\rho \mapsto g_\rho$ is smooth
- Let $T : \mathcal{A} \rightarrow \mathcal{B}$ be a quantum operation, then

$$T(\mathcal{D}) \subseteq \mathcal{D}(\mathcal{B}), \quad T(\mathcal{T}_\rho) \subseteq \mathcal{T}_{T(\rho)}$$

- Monotone metric if

$$\lambda_\rho(x, x) \geq \lambda_{T(\rho)}(T(x), T(x)), \quad x \in \mathcal{T}_\rho, \rho \in \mathcal{D},$$

A **quantum Fisher information** is a monotone Riemannian metric on \mathcal{D} .

Characterization of the quantum Fisher information

A function $f : \mathbb{R}^+ \rightarrow \mathbb{R}$ is **operator monotone** if $f(A) \leq f(B)$ whenever $0 \leq A \leq B$, $A, B \in B(\mathcal{K})$, $\dim(\mathcal{K}) < \infty$.

Theorem

(Petz (1996)) A Riemannian metric g on \mathcal{D} is monotone if and only if

$$g_\rho(x, y) = \text{Tr } x J_\rho^{-1}(y), \quad x, y \in \mathcal{T}_\rho$$

where $J_\rho = f(L_\rho R_\rho^{-1})R_\rho$ and $f : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ is an operator monotone function, such that $f(t) = tf(t^{-1})$.

We denote the metric corresponding to f by g^f .

Examples

- The SLD (Bures) metric: $f(t) = \frac{1+t}{2}$

$$g_{\rho}^f(x, y) = 2\text{Tr} (L_{\rho} + R_{\rho})^{-1}(x)y$$

- The BKM metric: $f(t) = \frac{t-1}{\log t}$

$$g_{\rho}^f(x, y) = \int_0^{\infty} \text{Tr} x(\rho + s)^{-1}y(\rho + s)^{-1} ds$$

- The RLD metric: $f(t) = \frac{2t}{1+t}$

$$g_{\rho}^f(x, y) = \frac{1}{2}\text{Tr} x(L_{\rho}^{-1} + R_{\rho}^{-1})(y)$$

Integral representation (Hansen (2005))

Let $f : \mathbb{R}^+ \rightarrow \mathbb{R}^+$ be operator monotone. Then

$$\frac{1}{f(t)} = \int_0^\infty \frac{1}{s+t} d\nu_f(s)$$

where ν_f is a positive Borel measure with support in $[0, \infty)$, such that $\int (1+s^2)^{-1} d\nu_f(s) < \infty$, $\int s(1+s^2)^{-1} d\nu_f + f(s) < \infty$.

It follows that

$$g_\rho^f(x, y) = \int_0^\infty \text{Tr } x(sR_\rho + L_\rho)^{-1}(y) d\nu_f(s)$$

Quantum χ^2 -divergence

The quantum χ^2 -divergence is given by

$$\chi_{1/f}^2(\sigma, \rho) = g_\rho^f(\sigma - \rho, \sigma - \rho)$$

where g^f is a quantum Fisher information. By monotonicity of g^f ,

$$\chi_{1/f}^2(\sigma, \rho) \geq \chi_{1/f}^2(T(\sigma), T(\rho))$$

for a quantum operation T .

Fisher information and reversibility

Denote by $\text{Lin}_\rho(\mathcal{S})$ the real linear span of $\mathcal{S} - \rho$. Then $\text{Lin}_\rho(\mathcal{S}) \subseteq \mathcal{T}_\rho$.

Theorem

Under the conditions 1. and 2., TFAE:

- (i) T is reversible for \mathcal{S} .
- (ii) $g_\rho^f(x, x) = g_{T(\rho)}^f(T(x), T(x))$ for all $x \in \text{Lin}_\rho(\mathcal{S})$ and all monotone metrics g^f .
- (iii) $\chi_{1/f}^2(\sigma, \rho) = \chi_{1/f}^2(T(\sigma), T(\rho))$ for all χ^2 divergences.
- (iv) Equality in (ii) holds for some f with $|\text{supp } \nu_f| \geq \dim(\mathcal{H})^2 + \dim(\mathcal{K})^2$.
- (v) Equality in (iii) holds for some f as in (iv).

Some open problems

- Support conditions for μ_f and ν_f : can be relaxed?
 - Cannot be removed: the supports must contain more than 1 point.
- Can we have full equivalence for L_1 -distance?
 - Up to now, only for extended family $\tilde{\mathcal{S}}$.
- Characterization by fidelity?
- Infinite dimensions: \mathcal{A} a von Neumann algebra?
 - Some results for f divergences, relative entropy (can be infinite!), S_α (Petz, Petz & AJ)
 - Characterization by commutant Radon-Nikodym derivative holds.
 - Factorization holds for type I algebras.
 - L_1 -distance?? QCB, QHB- ??